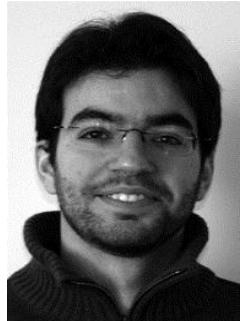


Smaller is different and more

Antonio M. García-García
EPSRC Career Acceleration
Cavendish Laboratory, Cambridge University

<http://www.tcm.phy.cam.ac.uk/~amg73/>



Pedro Ribeiro
Dresden



Santos & Way
Santa Barbara



Sangita Bose
Bombay

PRB, 86, 064526 (2012)
PRL 108, 097004 (2012)
PRB 84,104525 (2011)
Editor's Suggestion
Arxiv:1212.6779



Tezuka
Kyoto



Altshuler
Columbia

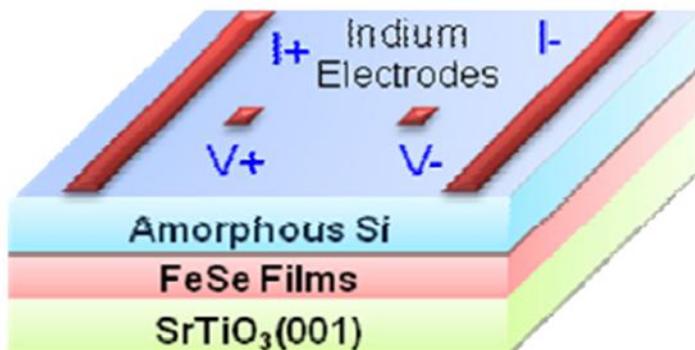
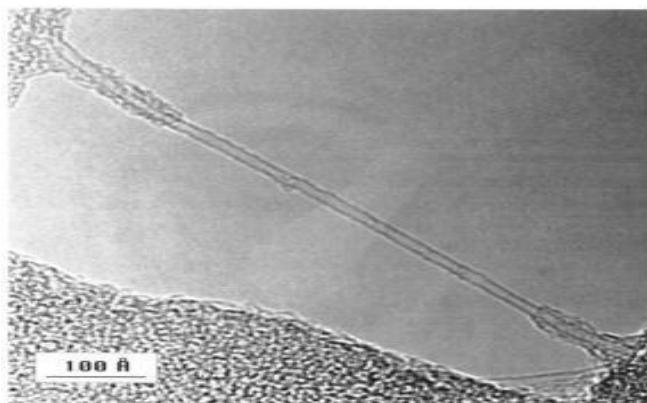
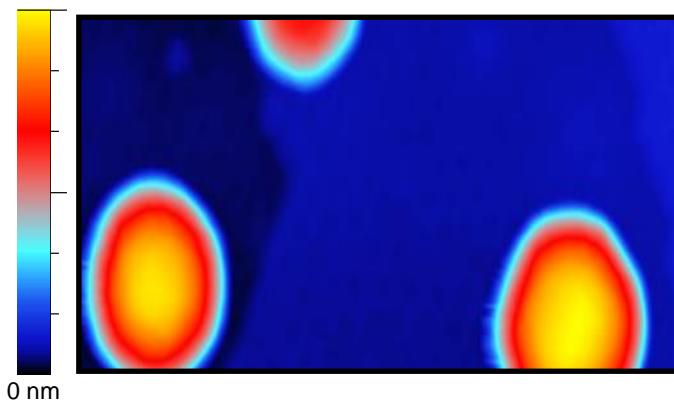


Klaus Kern
Stuttgart



Lobos
Maryland

7 nm



Why?

Boring?

Mermin-Wegner theorem

Low dimensions

=

No superconductivity

Exciting?

Beauty

Meso+Super
Quantum Control
Enhancement

Enhancement?

How to enhance
SC substantially?

with control

Mechanism of SC
in cuprates?

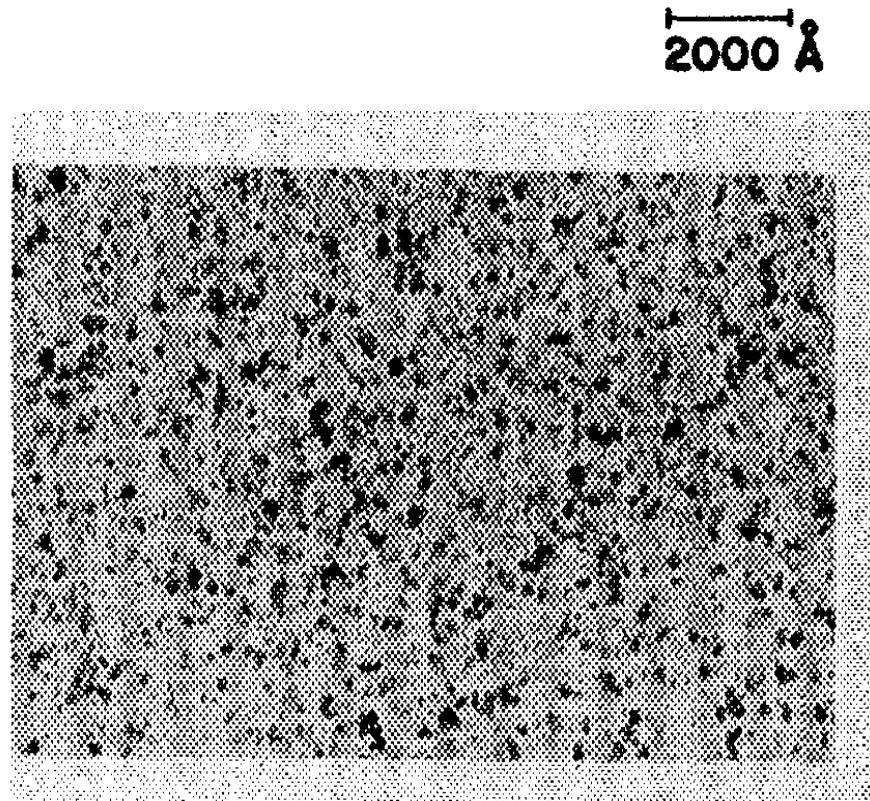
\$10⁶
Question



\$10
Question

Thin Films? JJ array?

Metal	T_c (°K)	T_c/T_{c0}	d (Å)	ρ_0
Al	3.0	2.6	40	0.19
Ga	7.2	6.5	...	0.20
Sn	4.1	1.1	110	0.31
In	3.7	1.1	110	0.36
Pb	7.2	1.0	...	0.53

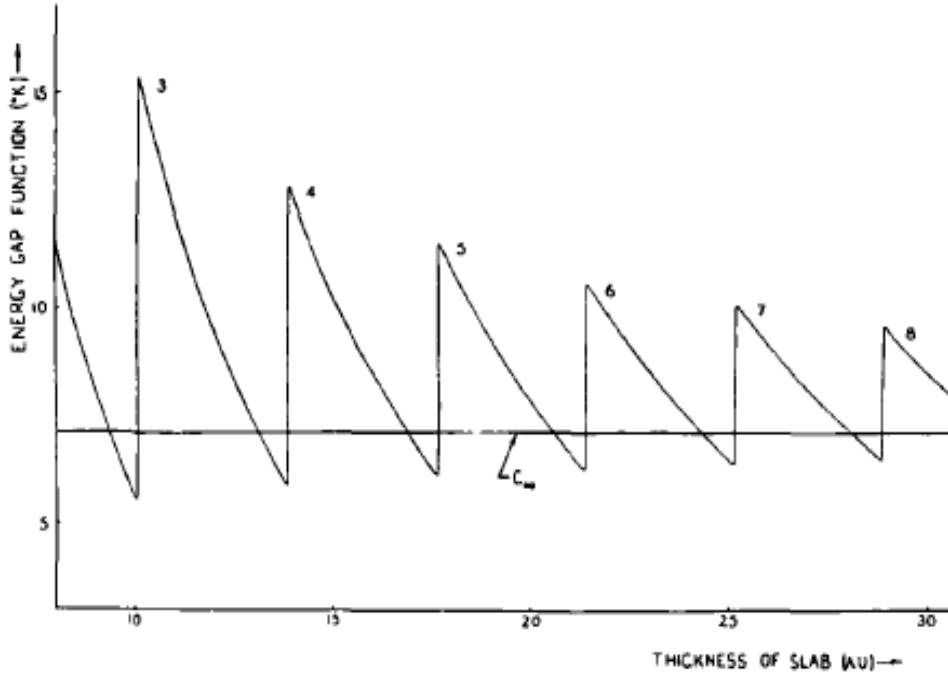


Abeles, Cohen, Cullen, Phys. Rev. Lett., 17, 632 (1966)

Crow, Parks, Douglass, Jensen, Giaver, Zeller....

A.M. Goldman, Dynes, Tinkham...

Thin Films



Single grains

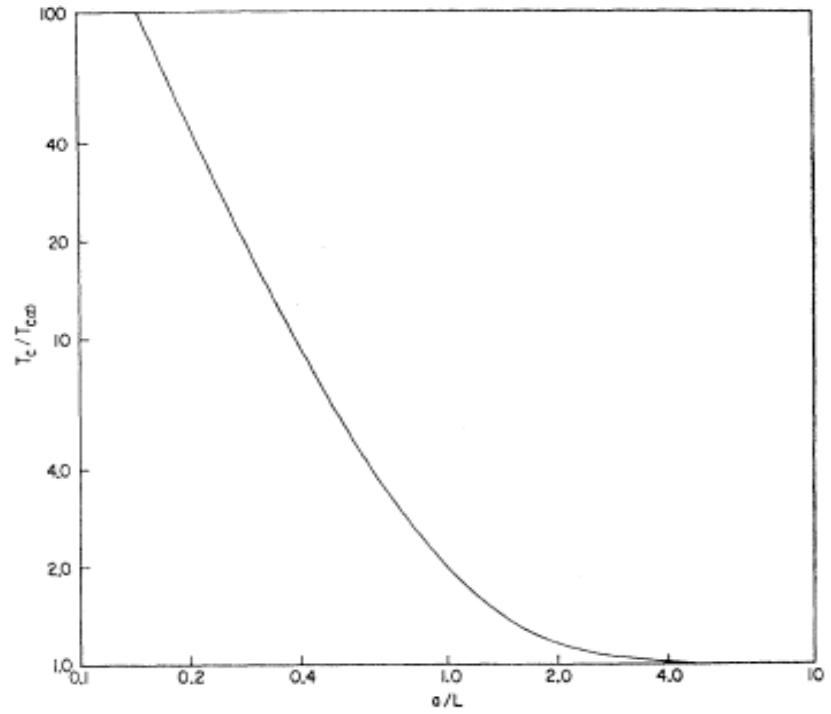


FIG. 1. $(T_c/T_{c\infty})$ versus (a/L) (see Ref. 17).

Shape Resonances

Blatt, Thompson
Phys. Lett. 5, 6 (1963)

Shell Effects

Parmenter, Phys. Rev. 166,
392 (1967)

Thinner

Smoothen

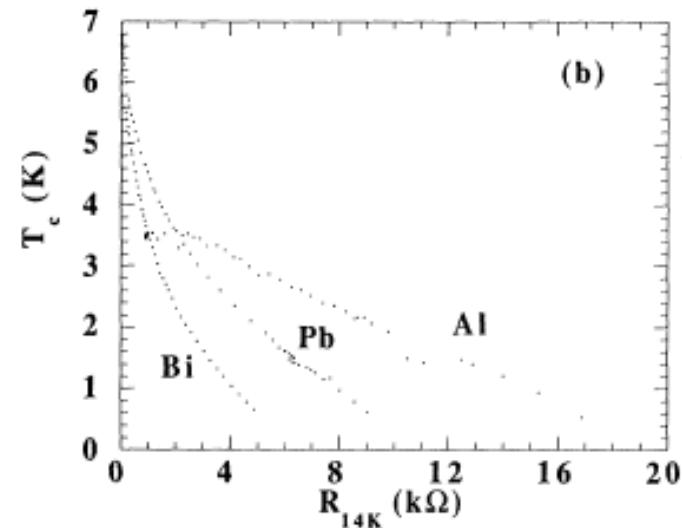
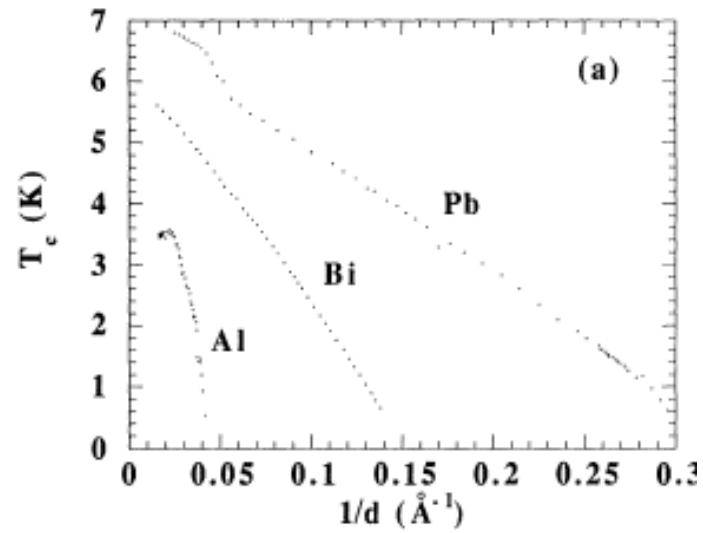
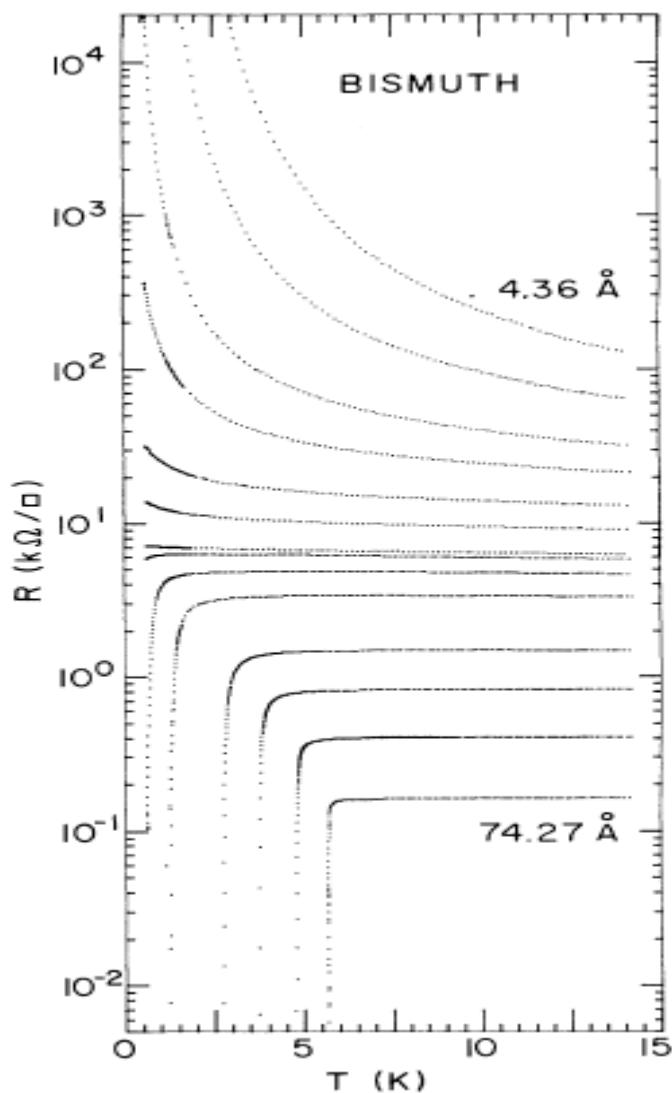
Disordered

BKT

Transition

$R_N > R_q$

Vortices
unbinding

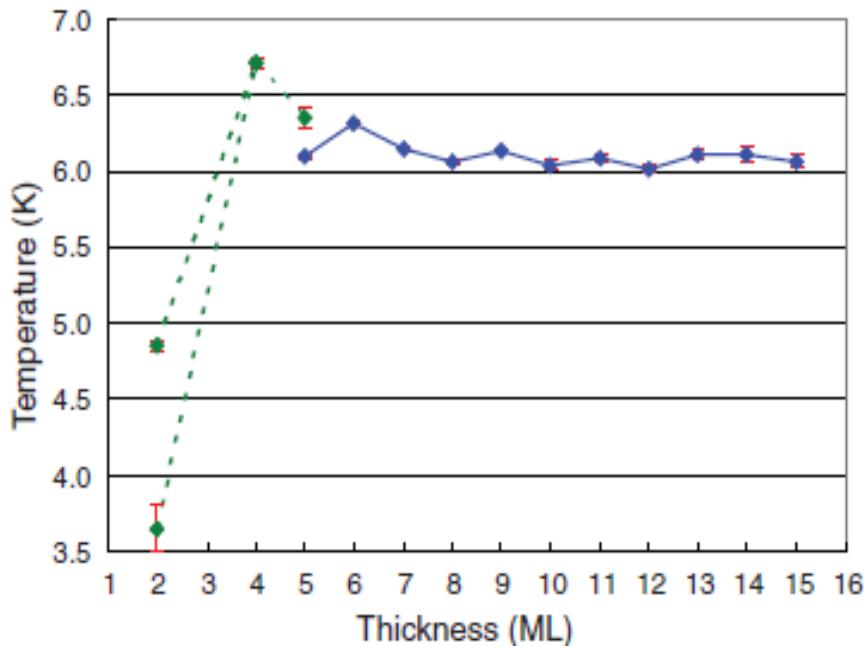


A.M. Goldman et al.

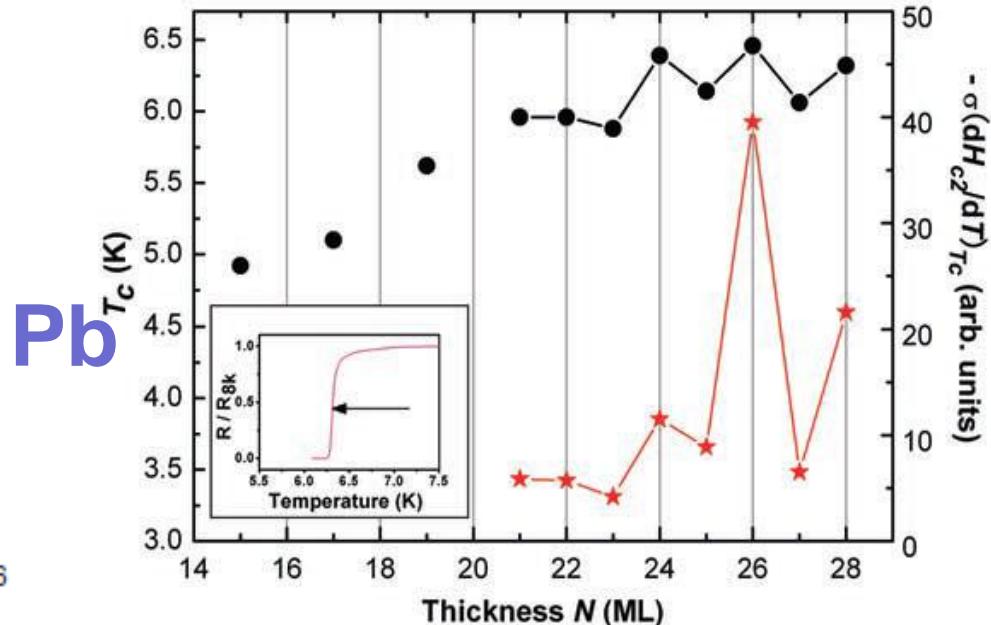
PRL 62 2180 (1989)
PRB 47 5931 (1993)

Recent

Atomic scale control



Pb

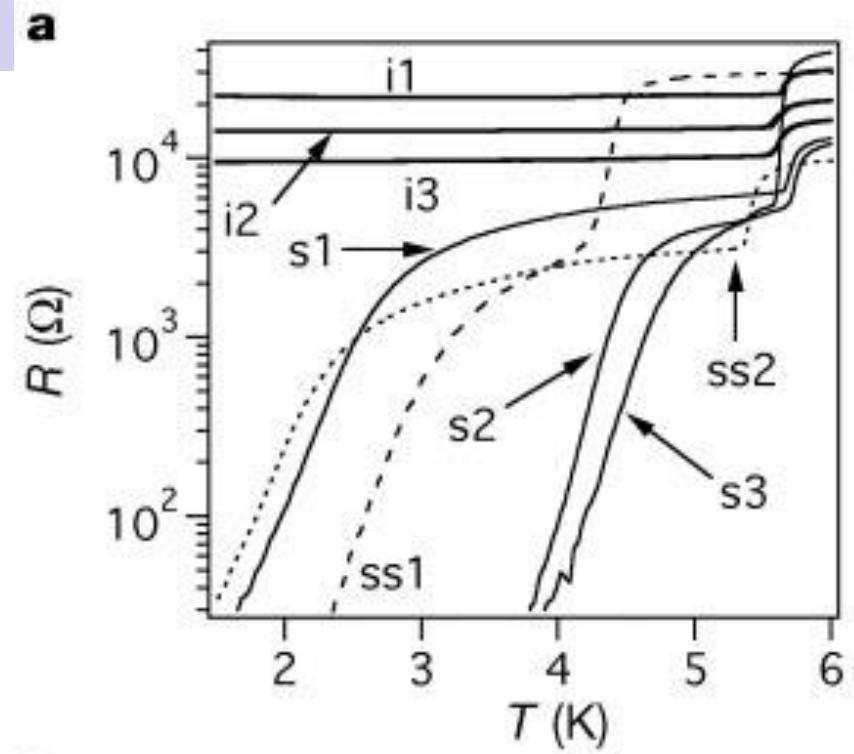
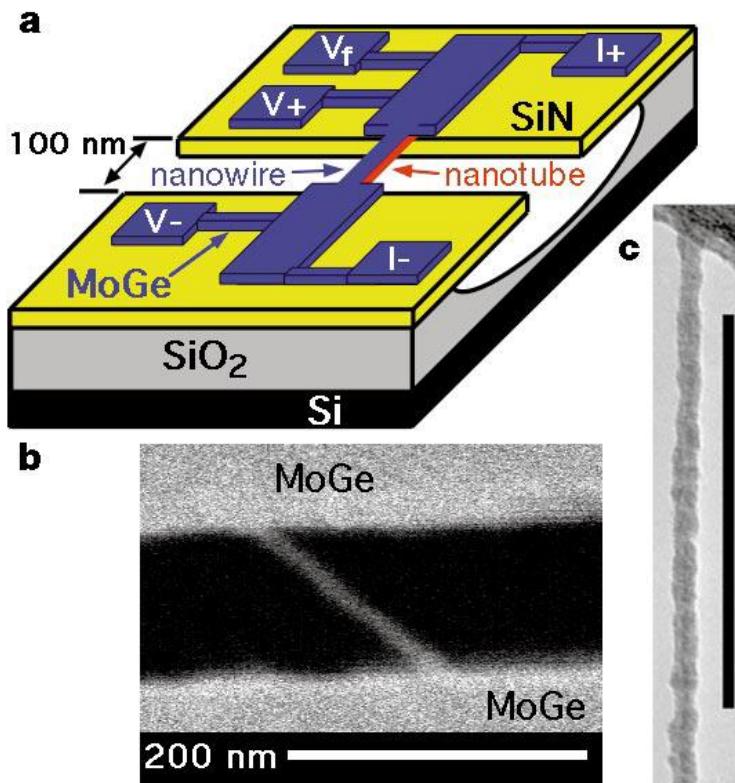


Shih et al., Science 324, 1314
(2009)

Xue et al., Science 306, 1915 (2004)

Xue et al., Nat Phys, 6 (2010), 104.

Nanowires $R \ll \xi$



Tinkham et al.
Nature 404, 971 (1990)

Superconductor
Insulator
transition

$$|\Delta(\mathbf{r}, t)| e^{i\theta(r,t)}$$

Fluctuation

$$\Delta(r_0, t_0) \approx 0$$

Finite
Resistance

Phase-slips

$$\theta \approx 0 \rightarrow 2\pi$$

$$R \propto e^{-S_{inst}}$$

Thermal

Langer & Ambegaokar,
PR. 164, 498 (1967).
McCumber & Halperin
PRB 1, 1054 (1970).

Quantum

Zaikin, A. D., Golubev, et al,
PRL 78, 1552 (1997).

Instantons

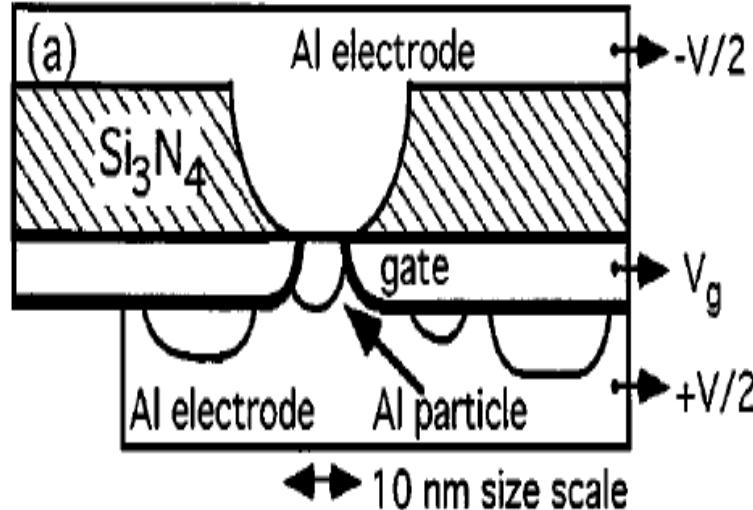
Small
Fluctuations

Coulomb-Gas

BKT transition

Quasi LRO

Ultrasmall Grains



Tinkham et al., PRL 74, 3244 (1995)

Richardson

It's exact. I did it
20 years ago

$T = 0$

$\delta/\Delta_0 \sim 1$

von Delft, Braun, Larkin, Sierra, Dukelsky,
Yuzbashyan, Matveev, Smith, Ambegaokar



Supercon-
ductivity?

1959

BCS fine until $\delta/\Delta_0 \sim 1/2$

BCS sharp transition

Richardson no transition

Richardson's equations

Von Delft, Yuzbashyan,
Dukelsky, Marsiglio, Braun
Sierra, Ambegaokar

$$-\frac{1}{\lambda d} + \sum_{j=1}^m' \frac{1}{E_i - E_j} = \frac{1}{2} \sum_{k=1}^n \frac{1}{E_i - \epsilon_k} \quad i = 1, \dots, m$$

Ground
state
energy

$$E = 2 \sum_{i=1}^m E_i + \sum_B \epsilon_B$$

Expansion
in δ/Δ_0

$$\Delta^b = 2\Delta_0 - d\sqrt{1 + \frac{\Delta_0^2}{D^2}} + \frac{d\Delta_0}{D} [1 + \phi(\lambda)]$$

Richardson, 1968, Yuzbashyan, Altshuler, 2005

$D \equiv E_D$
 $d \equiv \delta$

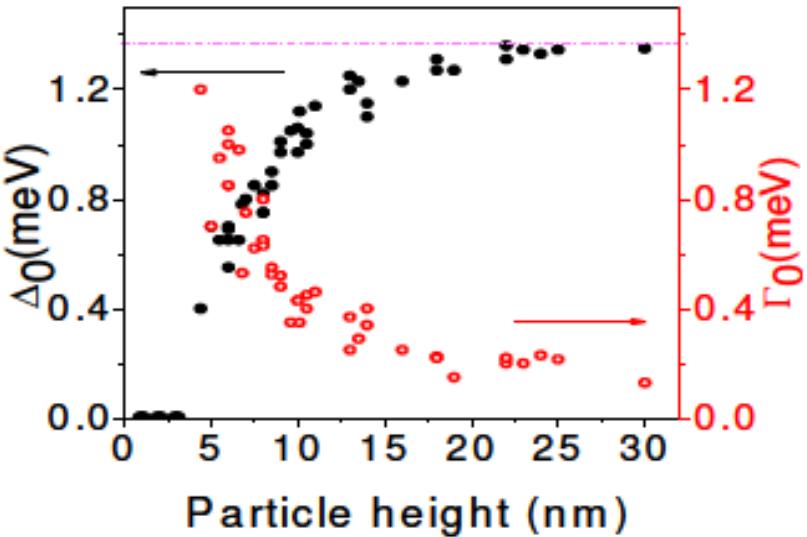
No enhancement



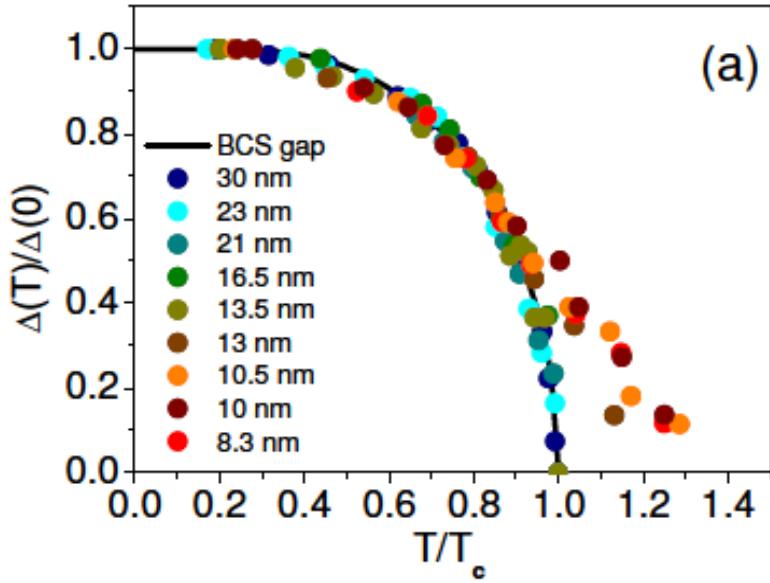
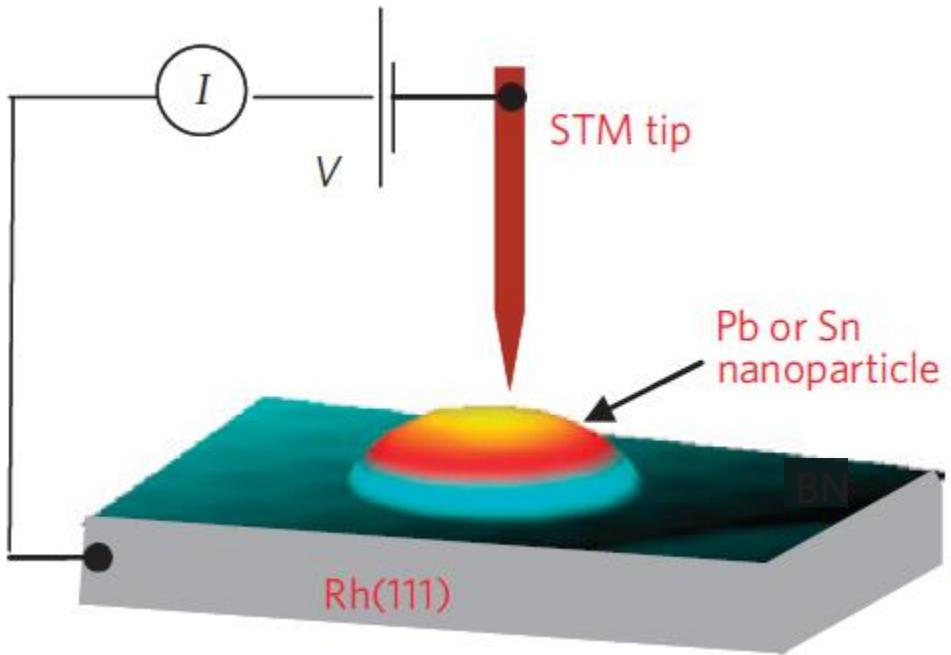
Kern
Stuttgart



Ribeiro
Dresden



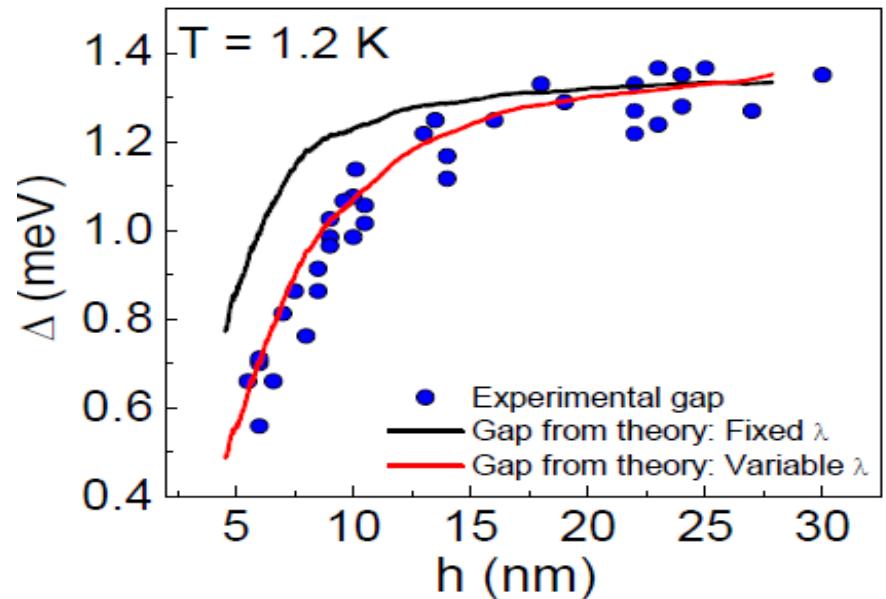
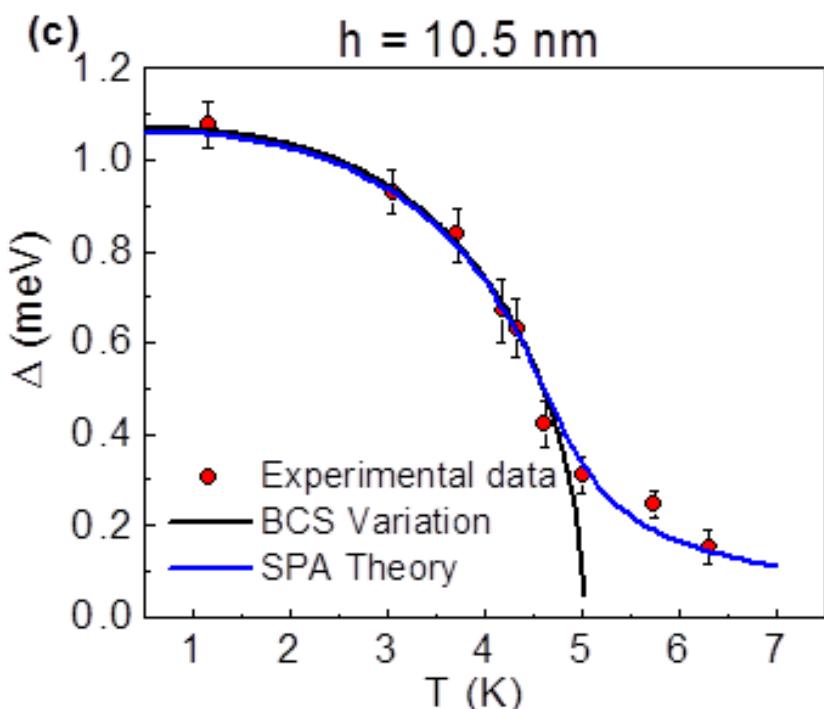
a



Quantum Fluctuations

Richardson's equations

and



Thermal Fluctuations

Static Path Approach

Scalapino, et al.

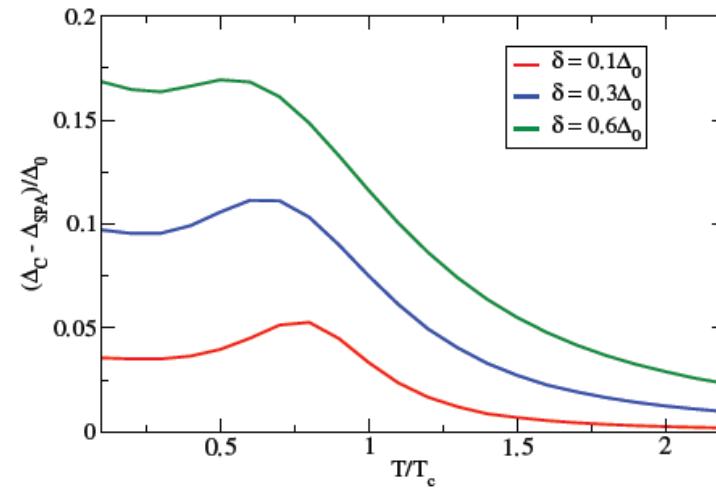
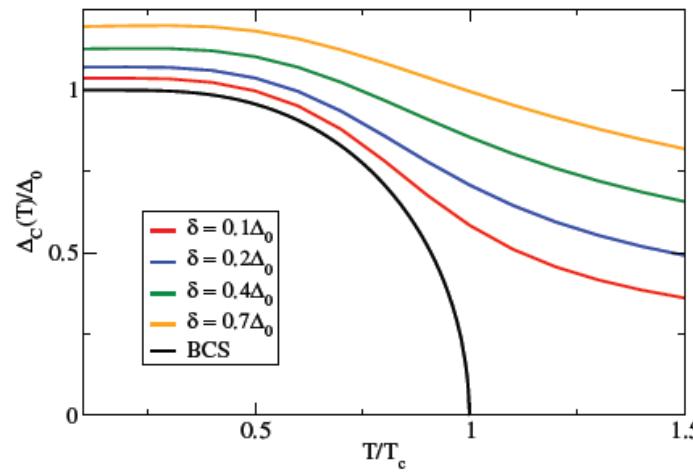
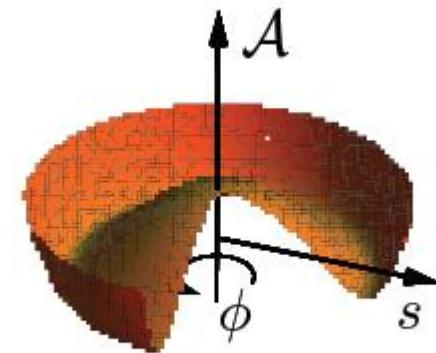
Divergences at intermediate T

Rossignoli and Canosa
Ann. of Phys. 275, 1, (1999)

$$\Delta(\tau) = s(\tau)e^{i\phi(\tau)}$$

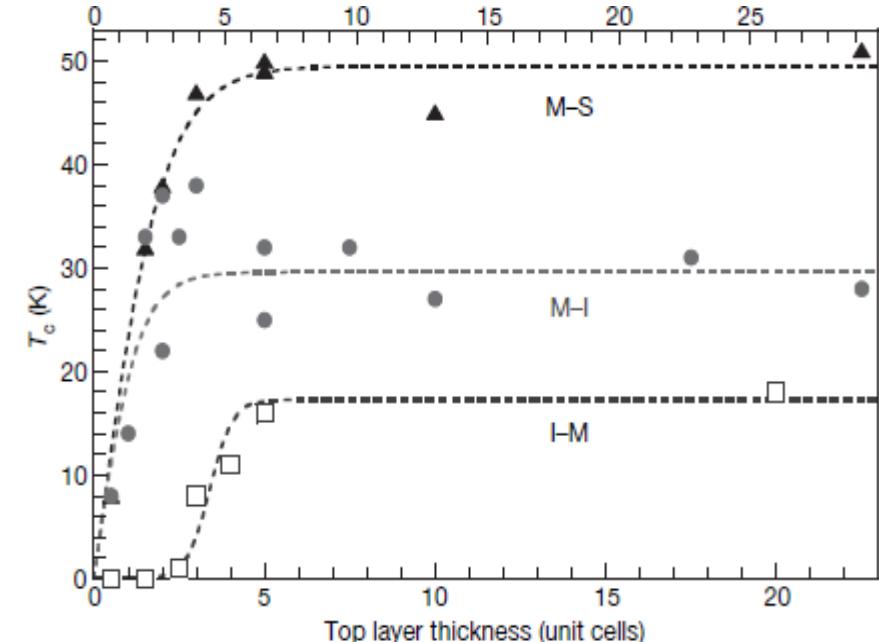
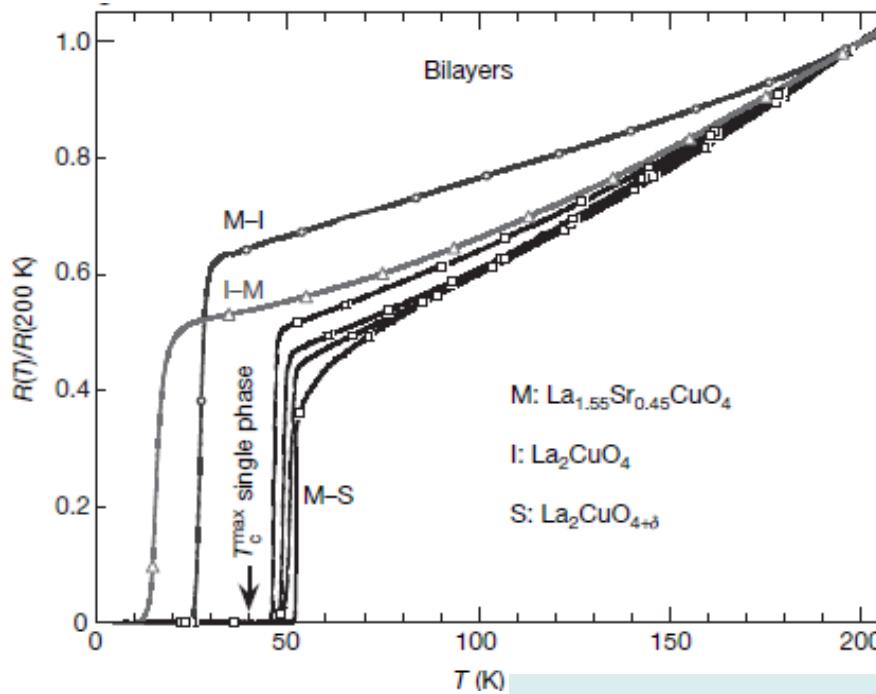
$$s^2(\tau) = s_0^2 + \delta s^2(\tau)$$

$$\phi(\tau) = \phi_0 + 2\pi M\tau/\beta + \delta\phi(\tau)$$



Is enhancement of
superconductivity
possible?

Cuprates high T_c Heterostructures

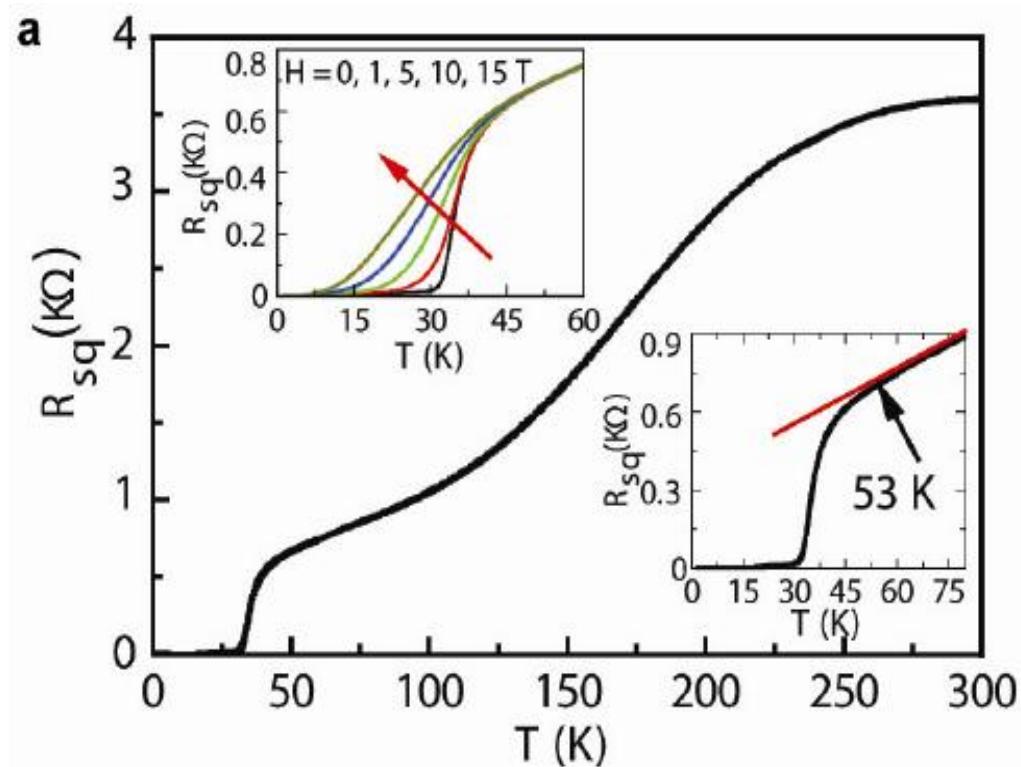
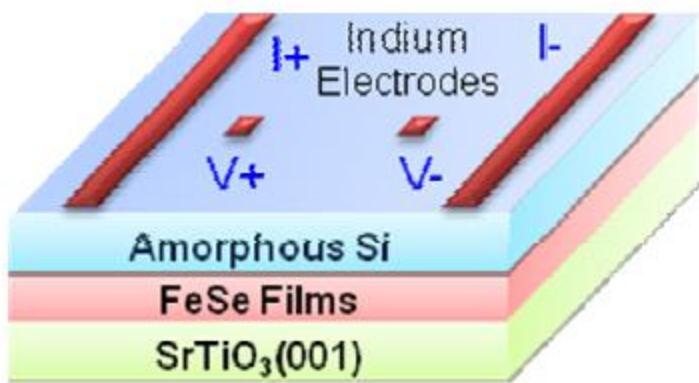


Bozovic et al., Nature 455, 782 (2008)

Higher T_c !!

Intrinsic inhomogeneities

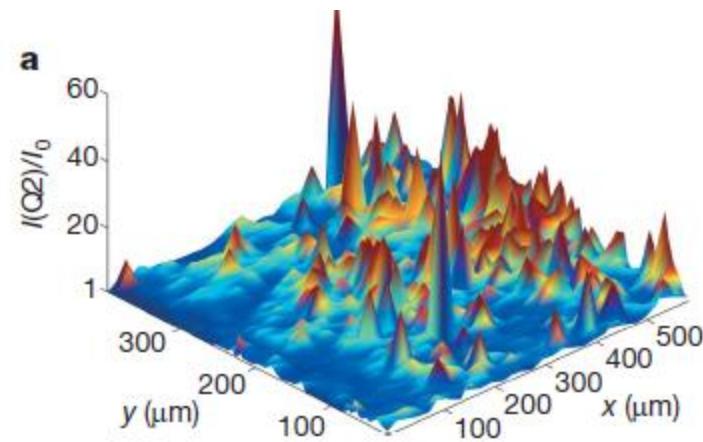
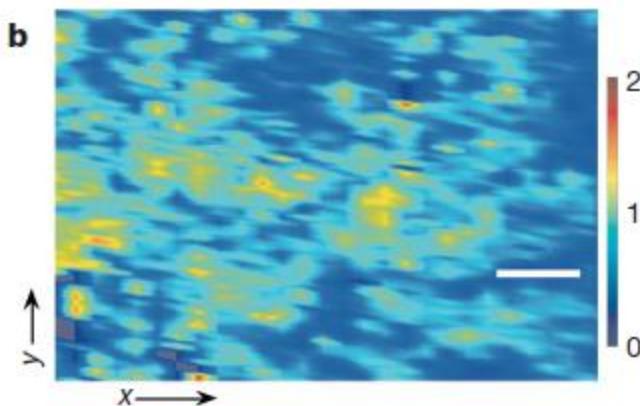
Iron Pnictides Heterostructures



Xue et al.
Nature Communications 3, 931 (2012)

Enhancement of T_c by disorder

Fractal distributions of dopants enhances SC in cuprates



Bianconi, et al., Nature 466, 841 (2010)

Inhomogeneities



Higher T_c

PRL 108, 017002 (2012)

$\Delta \gg \delta$

Grains

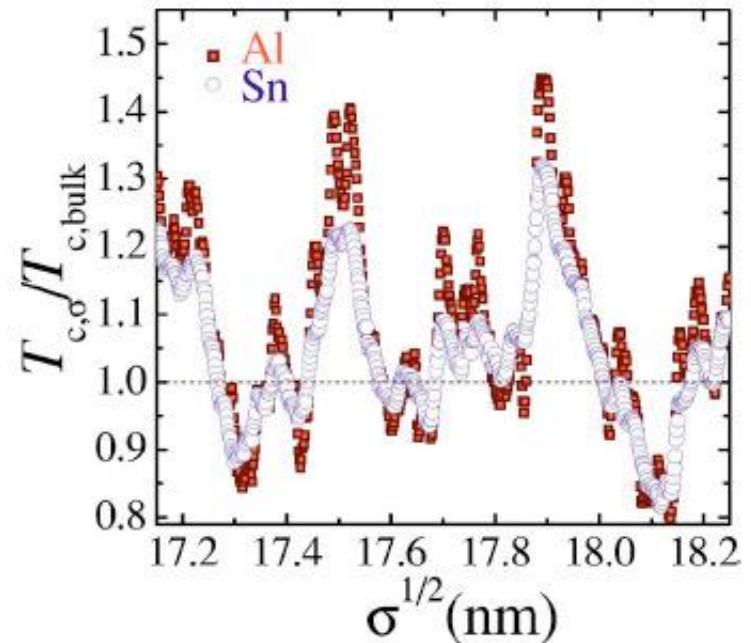
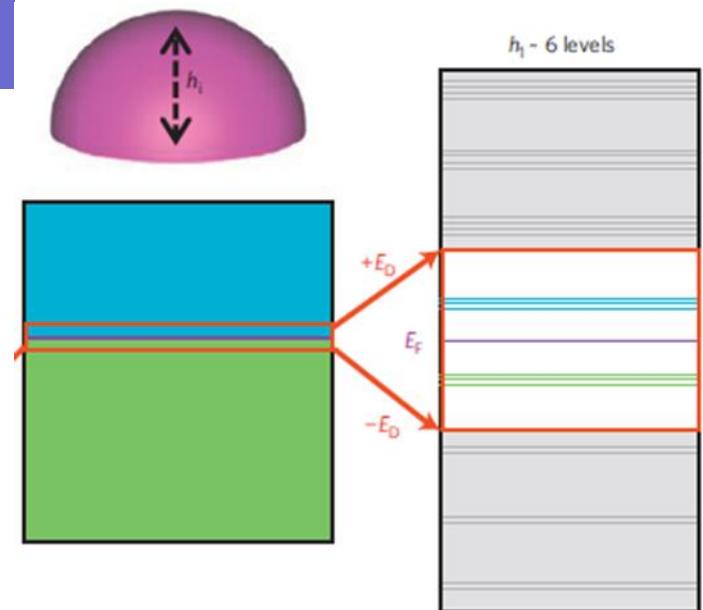
Heiselberg (2002): harmonic potentials, cold atom

Kresin, Ovchinnikov, Boyaci (2007) : Spherical, too high T_c

Peeters, et al, (2005-): BCS, BdG in a wire, cylinder..

Devreese (2006): Richardson equations in a box

Olofsson (2008): Estimation of fluctuations in BCS



Chaotic

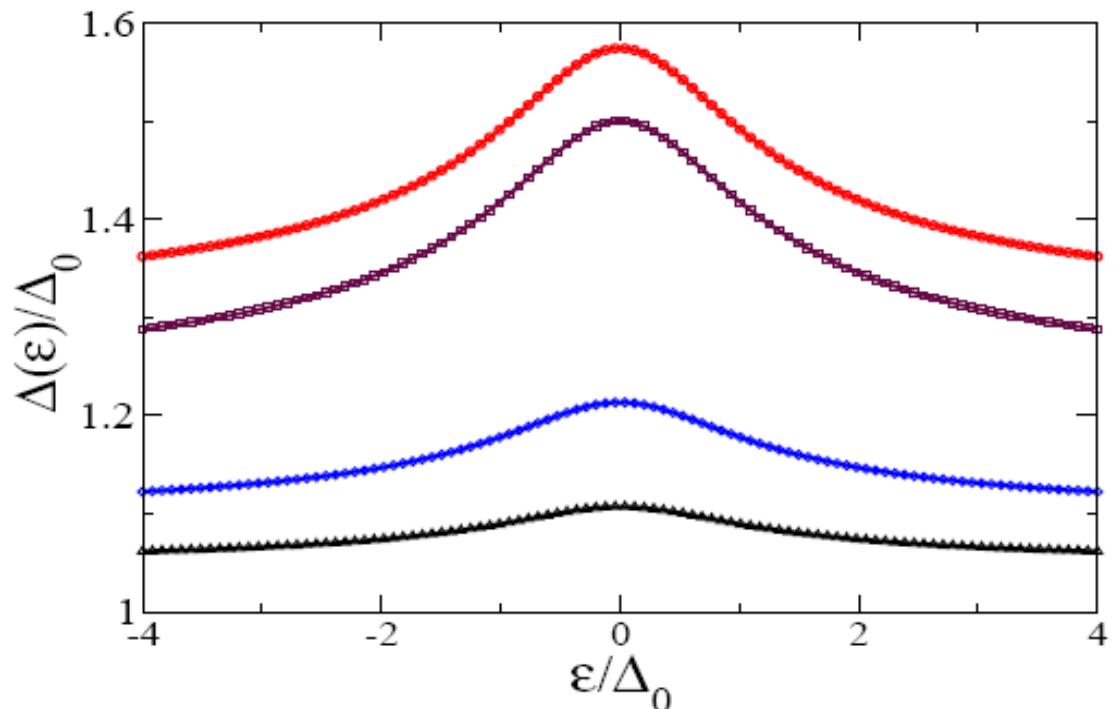
Al grain

$k_F = 17.5 \text{ nm}^{-1}$

$\Delta_0 = 0.24 \text{ mV}$

For $L < 9 \text{ nm}$ leading correction comes from density-density

AGG et al., PRL 100, 187001 (2008)
AGG et al., PRB 83, 014510 (2011)



$L = 6 \text{ nm}, \text{Dirichlet}, \delta/\Delta_0 = 0.67$

$L = 6 \text{ nm}, \text{Neumann}, \delta/\Delta_0 = 0.67$

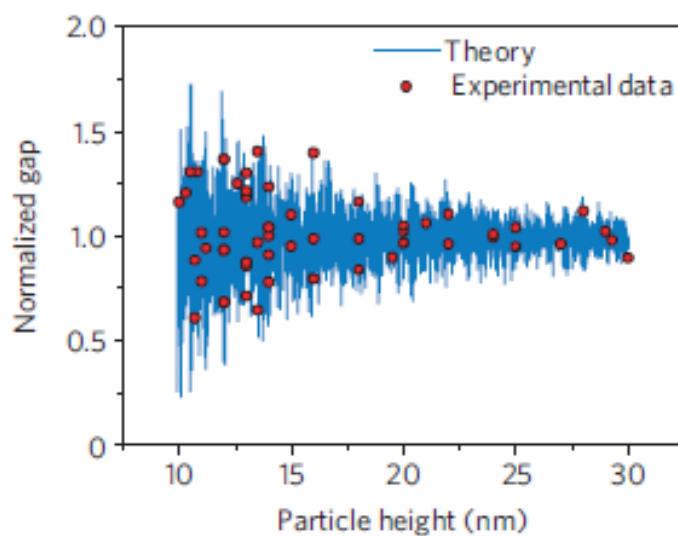
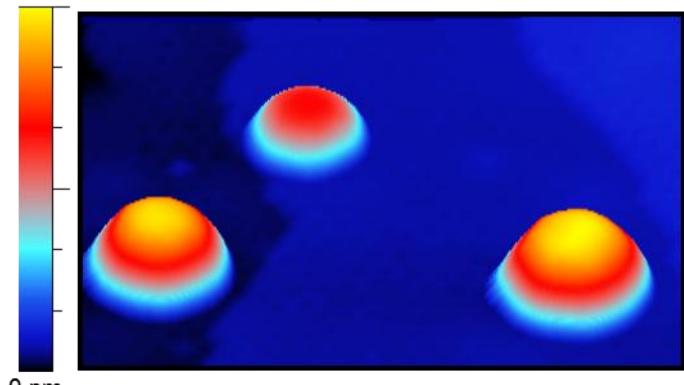
$L = 8 \text{ nm}, \text{Dirichlet}, \delta/\Delta_0 = 0.32$

$L = 10 \text{ nm}, \text{Dirichlet}, \delta/\Delta_0 = 0.08$



Single, Isolated Sn nano-grains

7 nm



Kern

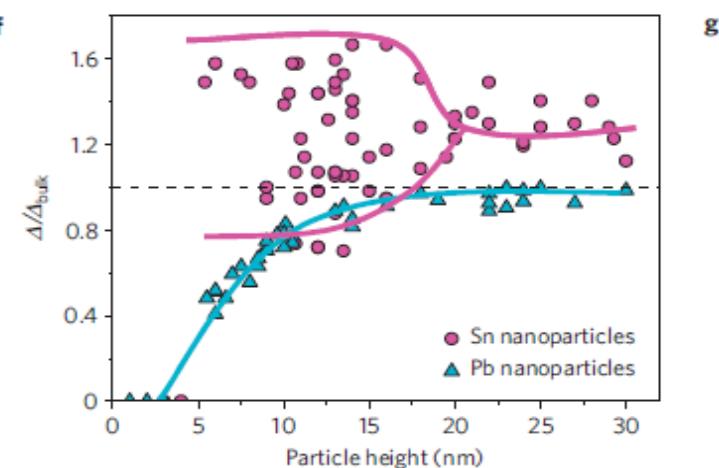


Bose

$R \sim 4\text{-}30\text{nm}$

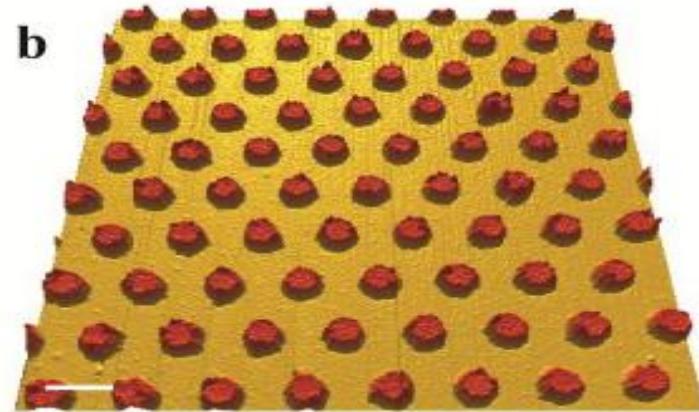
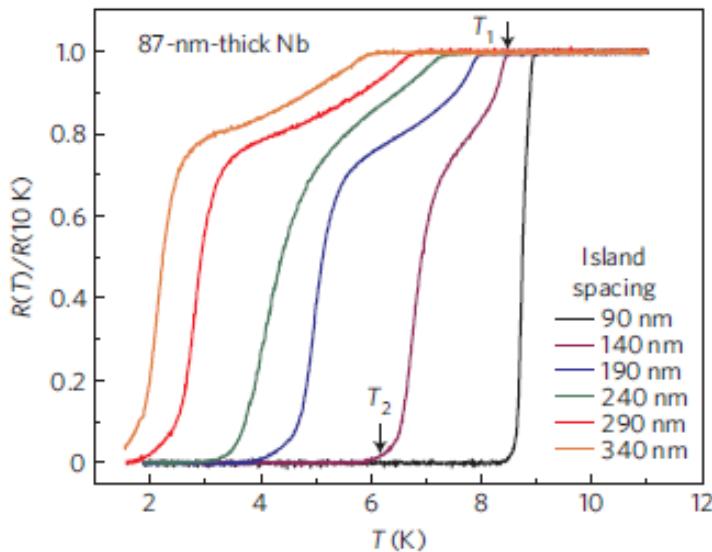
B closes gap

Almost hemispherical



True long-range order in nanograins?

Josephson junctions



Mason, Goldbart et al, Nature Physics 8 59
(2012)

James Mayoh and AGG, in preparation

$d = 1$

Quasi long-range order

1d + Dissipation = Long range order

Giamarchi, Blatter, Zaikin, Fisher, Lobos..



A. Lobos
Maryland

Why not
long-range
order in 1d?



M. Tezuka
Kyoto

Phase coherence in 1d by
power-law hopping

AGG et al., arXiv:1212.6779

1d Hubbard + power-law hopping

$$\begin{aligned}\mathcal{H} = & -\sum_{l \neq m, \sigma}^L \left(t_{lm} \hat{c}_{l,\sigma}^\dagger \hat{c}_{m,\sigma} + \text{H.c.} \right) - \mu \sum_{l=1, \sigma}^L \left(\hat{n}_{l,\sigma} - \frac{1}{2} \right) \\ & - |U| \sum_{l=1}^L \left(\hat{n}_{l,\uparrow} - \frac{1}{2} \right) \left(\hat{n}_{l,\downarrow} - \frac{1}{2} \right), \quad t_{lm} = t / |l - m|^\alpha\end{aligned}$$

$$|U|\gg t$$

$$\begin{aligned}\mathcal{H}_{\text{eff}} = & \frac{|U| L}{4} - \mu \sum_l (\hat{n}_l - 1) \\ & + \frac{4t^2}{|U|} \sum_{l \neq m} \left[\frac{(\hat{n}_l - 1) \hat{n}_m - \hat{\Delta}_l^\dagger \hat{\Delta}_m}{|l - m|^{2\alpha}} + \text{H.c.} \right]\end{aligned}$$

Previous



Bosonization

$$\rho(x) = \left[\rho_0 - \frac{\nabla\phi(x)}{\pi} \right] \sum_p e^{2ip(\pi\rho_0x - \phi(x))}$$

$$\Delta(x) = \rho_0 e^{-i\theta(x)} \sum_p e^{2ip(\pi\rho_0x - \phi(x))},$$

$$[\nabla\phi(x), \theta(y)] = i\pi\delta(x-y)$$

Phase

Density

$$\langle \Delta(x) \rangle = \langle \hat{c}_{x,\uparrow}^\dagger \hat{c}_{x,\downarrow}^\dagger \rangle \propto \langle e^{-i\theta(x)} \rangle \quad \delta\rho(x) \simeq -\nabla\phi(x)/\pi$$

$$\begin{aligned} \mathcal{H}_{\text{eff}} = & \int dx \left[\tilde{\mu} \frac{\nabla\phi(x)}{\pi} + \frac{uK}{2\pi} (\nabla\theta(x))^2 + \frac{u}{2\pi K} (\nabla\phi(x))^2 \right] \\ & - g \frac{\pi u \rho_0^2}{4Ka^{1-2\alpha}} \int_{|x-x'|>a} dx dx' \frac{\cos [\theta(x) - \theta(x')]}{|x-x'|^{2\alpha}}. \end{aligned}$$

Self-consistent harmonic approximation (SCHA)

$$S_0 = \frac{1}{2\beta L} \sum_{\mathbf{q}} g_0^{-1}(\mathbf{q}) \theta_{\mathbf{q}}^* \theta_{\mathbf{q}}$$

$$F_{\text{var}} = F_0 + T \langle S - S_0 \rangle_0$$

$$\begin{aligned} g_0^{-1}(\mathbf{q}) &= \frac{K}{\pi u} \omega_m^2 + \frac{uK}{\pi} k^2 + \frac{2\pi u \rho_0^2}{Ka^{2-2\alpha}} \int_a^L dr \frac{1 - \cos kr}{r^{2\alpha}} \\ &\quad \times \exp \left[-\frac{1}{\beta L} \sum_{\mathbf{q}'} (1 - \cos k'r) g_0(\mathbf{q}') \right] \end{aligned}$$

$$g_0^{-1}(\mathbf{q}) = \frac{K}{\pi u} \omega_m^2 + \frac{uK}{\pi} k^2 + \eta |k|^{2\alpha-1}$$

$\alpha < 3/2$

Long-range order

$\alpha > 3/2$

Quasi long-range order

$$\langle \Delta(x_i) \rangle \sim \langle e^{i\theta(x_i)} \rangle \neq 0$$

Phase slips suppressed

$d = 1$

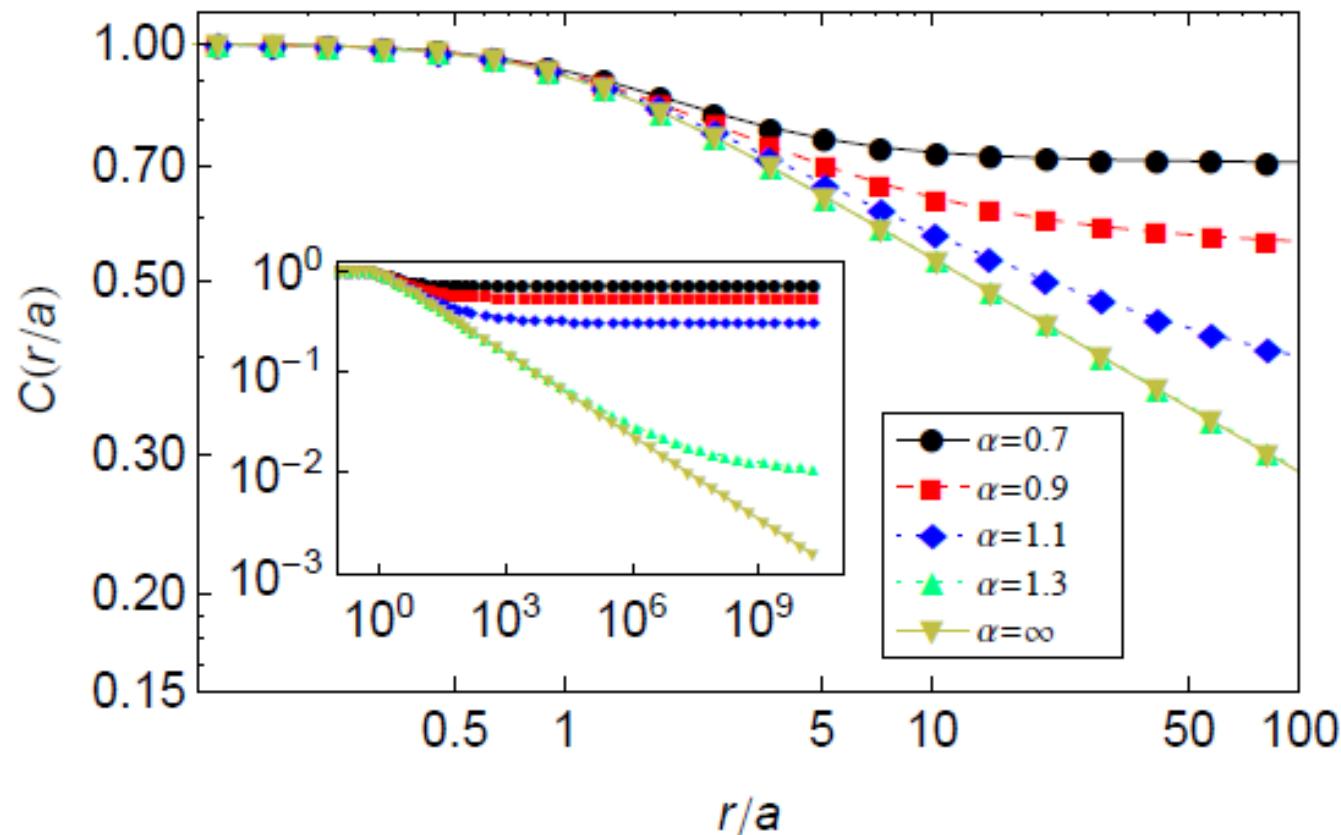
BUT

$$d_{eff} = 2/(2\alpha - 1)$$

$d_{eff} > 1$

Phase coherence

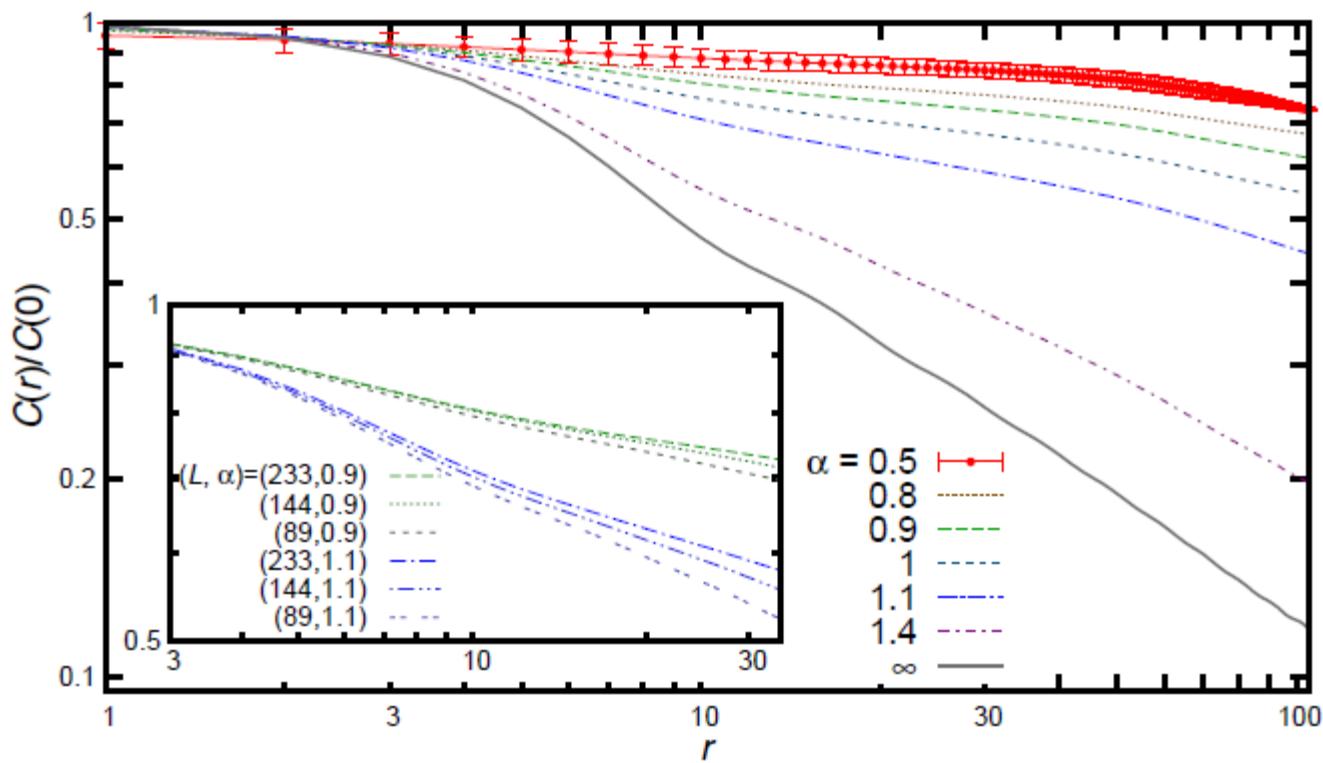
Bosonization:Correlations



$$C(r) = \langle e^{i\theta(r)} e^{-i\theta(0)} \rangle_0$$

$$C(r) \approx e^{-G(0)} \left[1 + A/r^{3/2-\alpha} + \mathcal{O}\left(1/r^{3-2\alpha}\right) \right]$$

Numerical: DMRG



$$C(r) \equiv \frac{1}{L - 2l_0 - r} \sum_{l=l_0+1}^{L-l_0-r} \langle \hat{\Delta}_{l+r} \hat{\Delta}_l^\dagger \rangle$$

$$\hat{\Delta}_l^\dagger \equiv \hat{c}_{l,\uparrow}^\dagger \hat{c}_{l,\downarrow}^\dagger$$

From SC to Quantum magnetism

$$\begin{aligned}\mathcal{H}_{\text{eff}} = & \frac{|U| L}{4} - \mu \sum_l (\hat{n}_l - 1) \\ & + \frac{4t^2}{|U|} \sum_{l \neq m} \left[(-1)^{(\hat{n}_l - 1)\hat{n}_m} \frac{\hat{\Delta}_l^\dagger \hat{\Delta}_m}{|l - m|^{2\alpha}} + \text{H.c.} \right]\end{aligned}$$

Anderson
Pseudo-spin
representation

$$\hat{n}_l \rightarrow \hat{S}_l^z + 1/2$$

$$\hat{\Delta}_l^\dagger \rightarrow \hat{S}_l^+$$

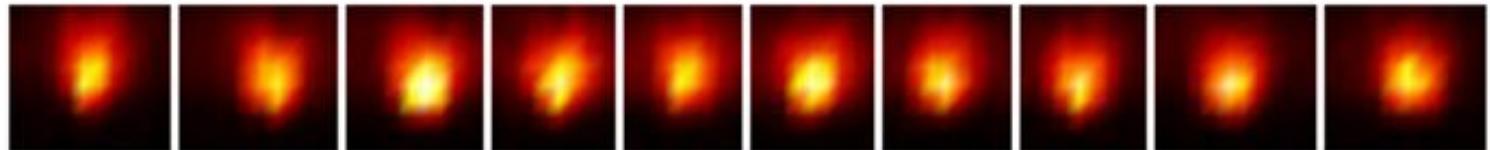
XXZ Spin
chain

Heisenberg
chain

Variable-Range Interactions in Trapped Ion Quantum Simulators

Islam, Monroe.. 1210.0142, Bollinger, Britton... Nature 484, 489 (2012)

$^{171}Yb^+$



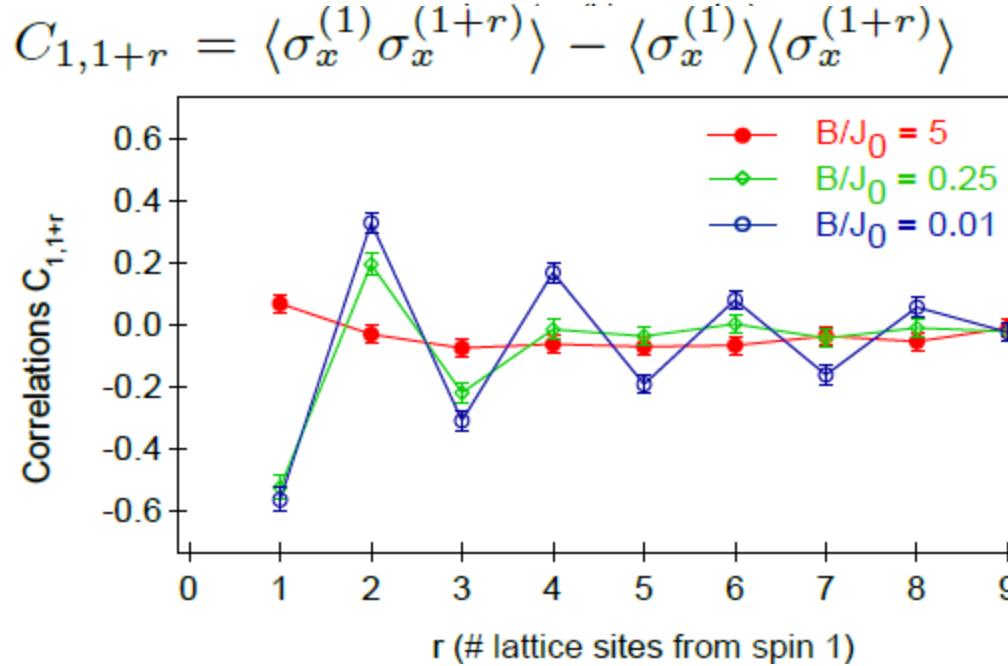
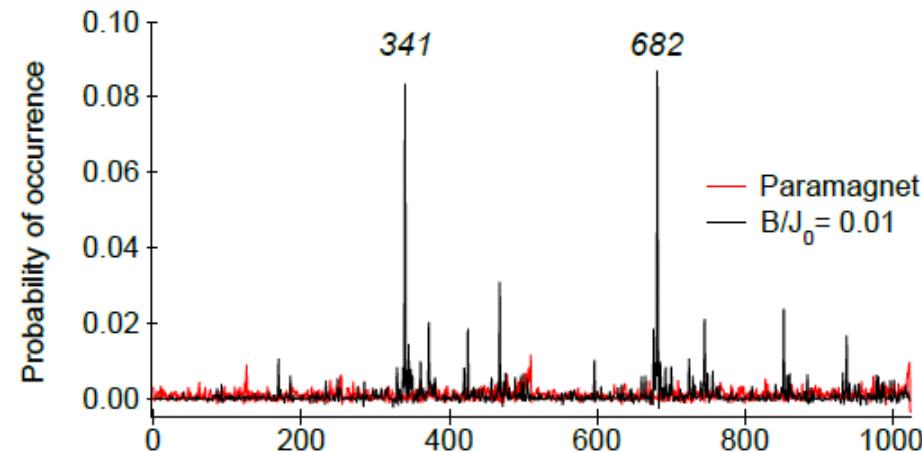
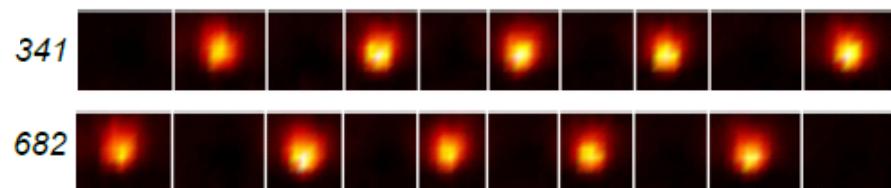
Raman transitions

$\pi/2$ -shift

Dipole forces

$$H = \sum_{j < i} J_{ij} \sigma_x^{(i)} \sigma_x^{(j)} - B \sum_i \sigma_y^{(i)} \quad J_{ij} \approx \frac{J_0}{|i - j|^\alpha}$$

$$0 < \alpha < 3$$



(anti)Ferromagnetic Transitions

Frustration

Spin Liquids

Quantum Magnetism

Nano meets high

Holographic grains

Finite Size + Strong interactions ?

Tough for even
conventional
superconductors



Way
Santa Barbara

Holographic
superconductivity in
confined geometries?



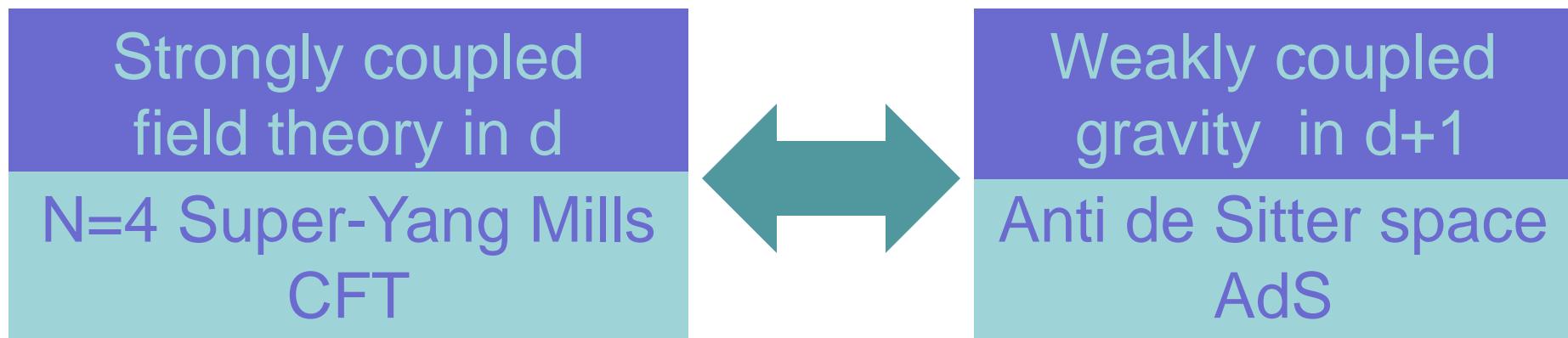
Santos
Santa Barbara

Holographic principle

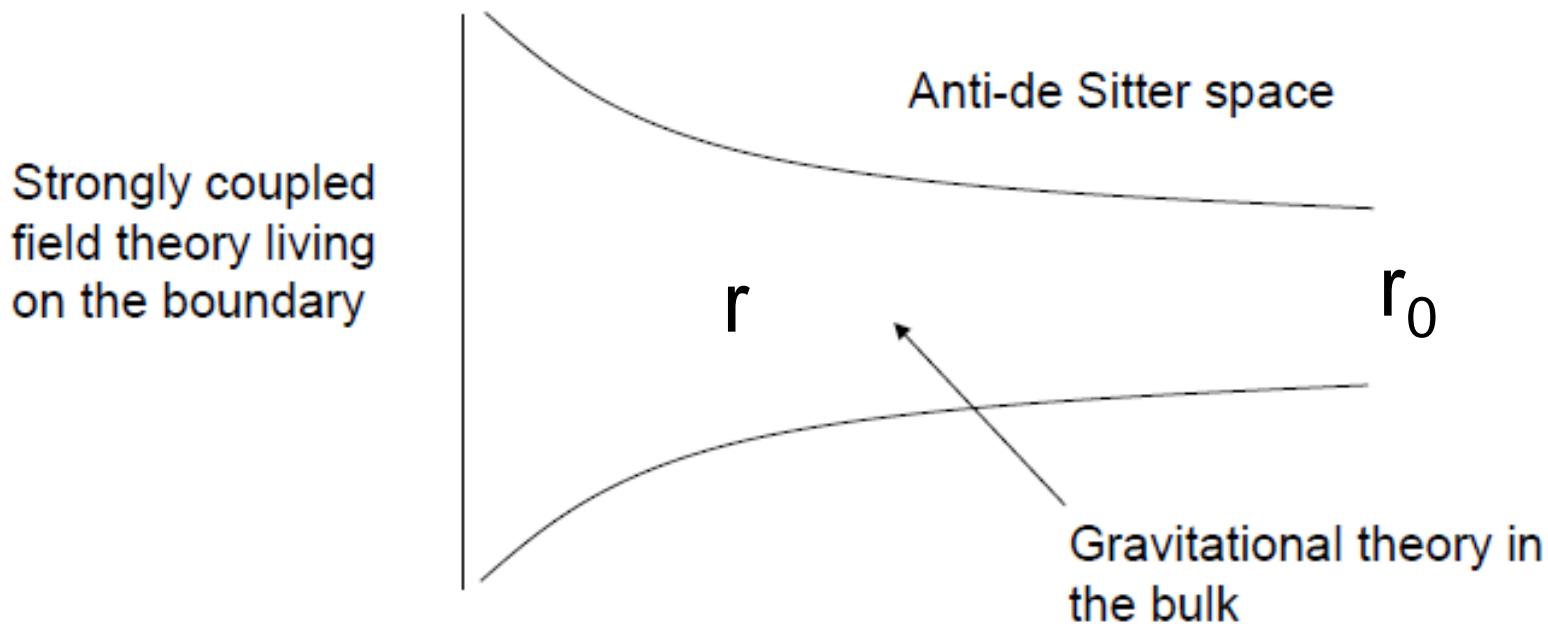
Maldacena's conjecture

AdS/CFT correspondence

t'Hooft, Susskind, Weinberg, Witten....



Extra dimension? Geometrization of Wilson RG



Holography beyond string theory

2003

QCD Quark gluon plasma
Gubser, Son

2008

Holographic superconductivity
Hartnoll, Herzog, Horowitz

2012

Quantum criticality, non-equilibrium..
Zaanen, Sachdev, Philips

Easy to compute in the
gravity dual

&

Detailed
dictionary

An answer looking for a question

$H = ?$

I do not know

I know
that

Complex scalar

Spontaneous breaking
 $U(1)$ at low T

Finite μ

Simplest dual gravity theory

$$S = \int d^4x \sqrt{-g} \left[R + \frac{6}{L^2} - \frac{1}{4} F_{\mu\nu} F^{\mu\nu} - |D\psi|^2 - m^2 |\psi|^2 \right]$$

$$D = \nabla - iqA$$

ψ ≡ complex scalar

Metric

$$\begin{aligned} ds^2 &= -f(r)dt^2 + \frac{dr^2}{f(r)} + r^2(dx^2 + dy^2) \\ &= -f(r)dt^2 + \frac{dr^2}{f(r)} + r^2(dR^2 + R^2d\theta^2) \\ f(r) &= \frac{r^2}{L^2} \left(1 - \frac{r_0^3}{r^3} \right) , \end{aligned}$$

Equations of motion:

$$\partial_r^2 |\psi| + \frac{1}{r^2 f} \partial_x^2 |\psi| + \left(\frac{f'}{f} + \frac{2}{r} \right) \partial_r |\psi| + \frac{1}{f} \left(\frac{A_t^2}{f} - m^2 \right) |\psi| = 0$$

$$\partial_r^2 A_t + \frac{1}{r^2 f} \partial_x^2 A_t + \frac{2}{r} \partial_r A_t - \frac{2|\psi|^2}{f} A_t = 0$$

Boundary conditions:

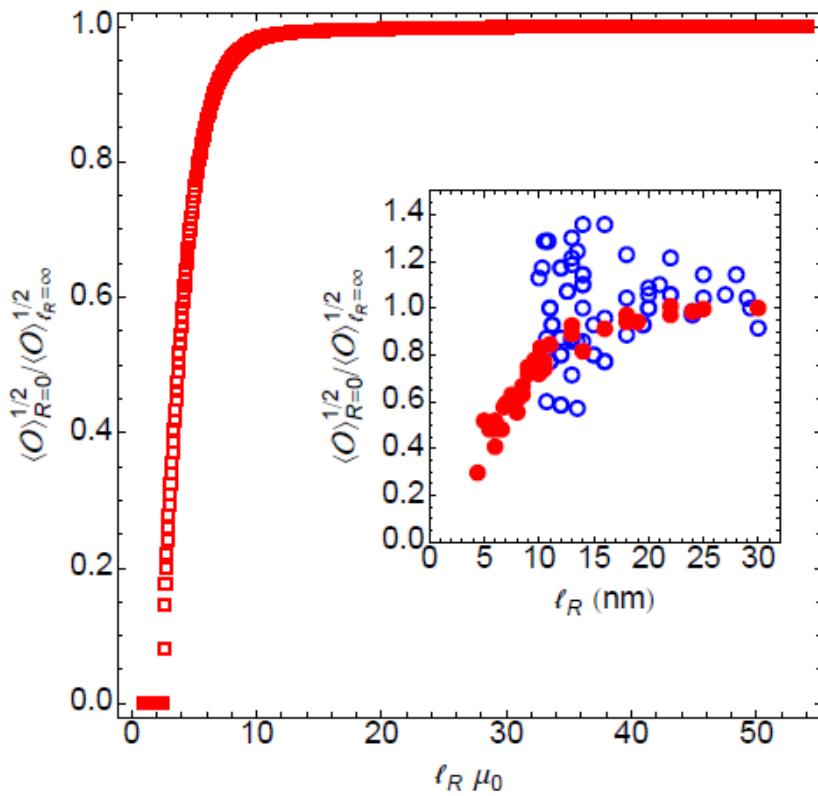
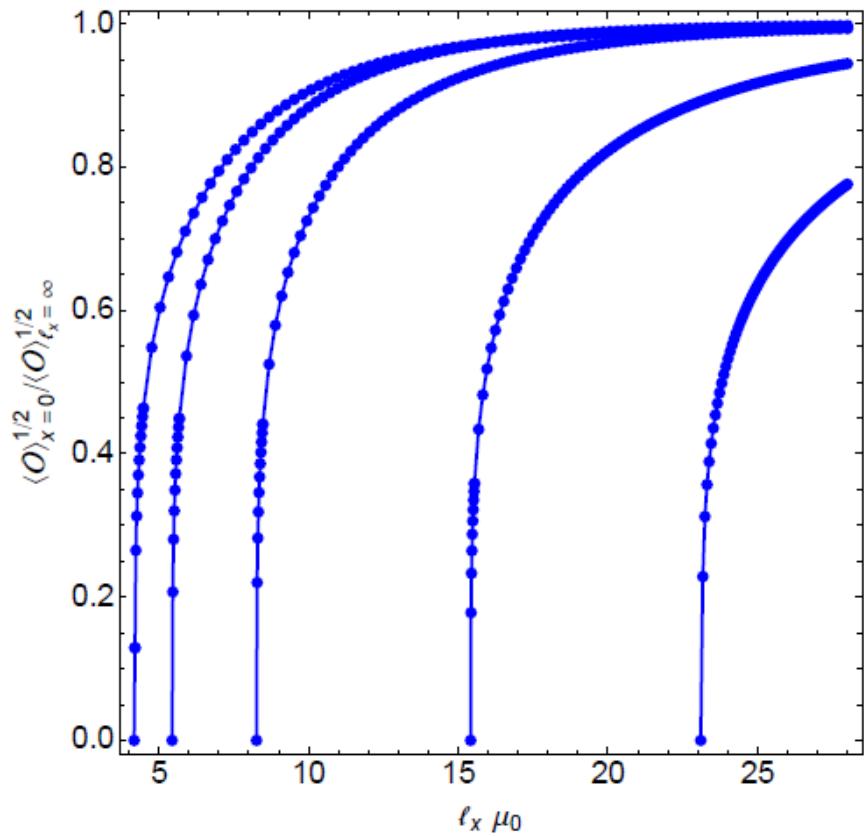
$$\begin{aligned} \mathbf{r} &= \mathbf{r}_0 & \mathbf{r} \rightarrow \infty & |\psi| = \frac{\psi^{(1)}}{r} + \frac{\psi^{(2)}}{r^2} + O\left(\frac{1}{r^3}\right) \\ \mathbf{A}_t &= \mathbf{0} & & A_t = \mu - \frac{\rho}{r} + O\left(\frac{1}{r^2}\right) \end{aligned}$$

How small?

$$\mu(x) = \mu_0 \left[\frac{1 - \epsilon + \epsilon \cosh\left(\frac{2x}{\sigma}\right) + \cosh\left(\frac{\ell_x}{\sigma}\right)}{\cosh\left(\frac{2x}{\sigma}\right) + \cosh\left(\frac{\ell_x}{\sigma}\right)} \right]$$

Dictionary:

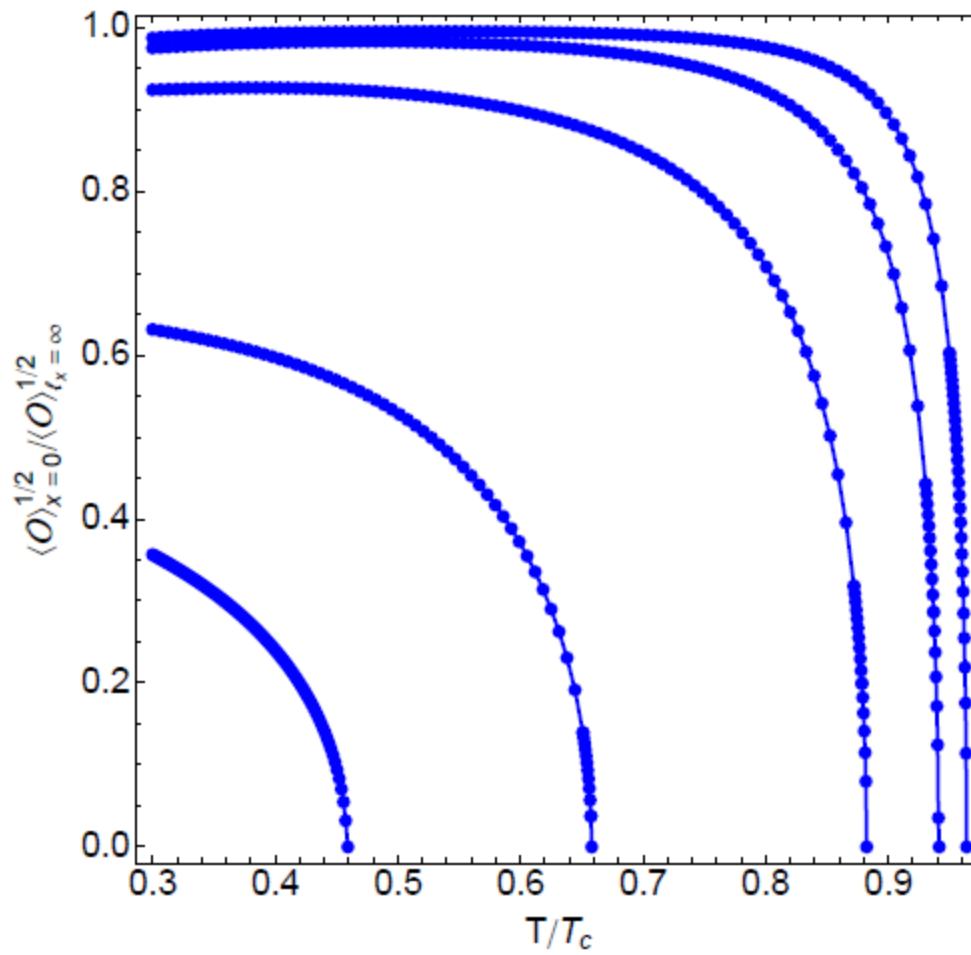
$$\langle \mathcal{O} \rangle = \sqrt{2} \psi^{(2)}$$



“Superconductivity” only for $|l| < l_c$

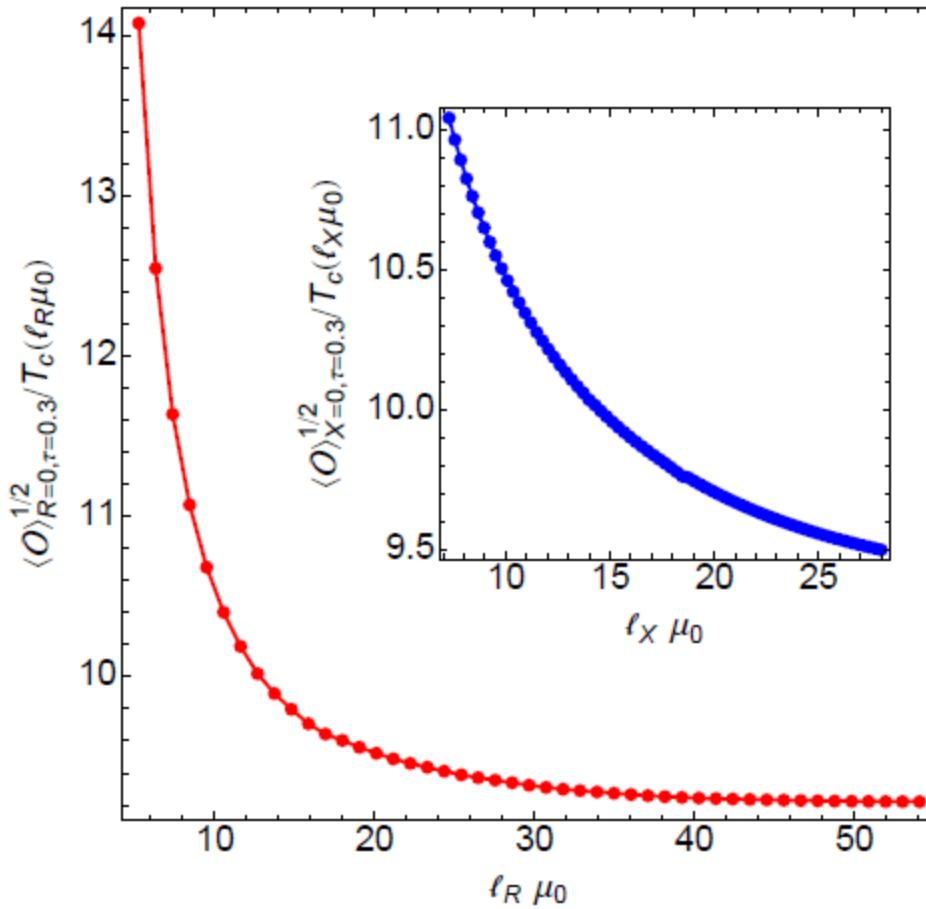
Mean field behavior

Fluctuations?



No thermal fluctuations

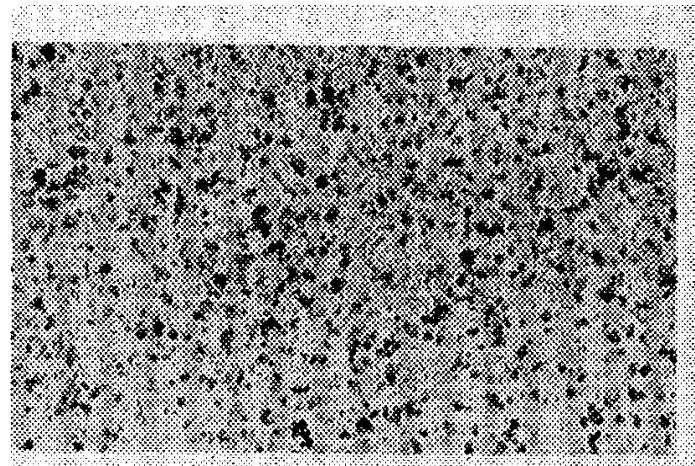
Large N artefact



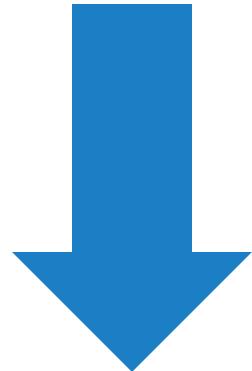
Interactions depend on system size!

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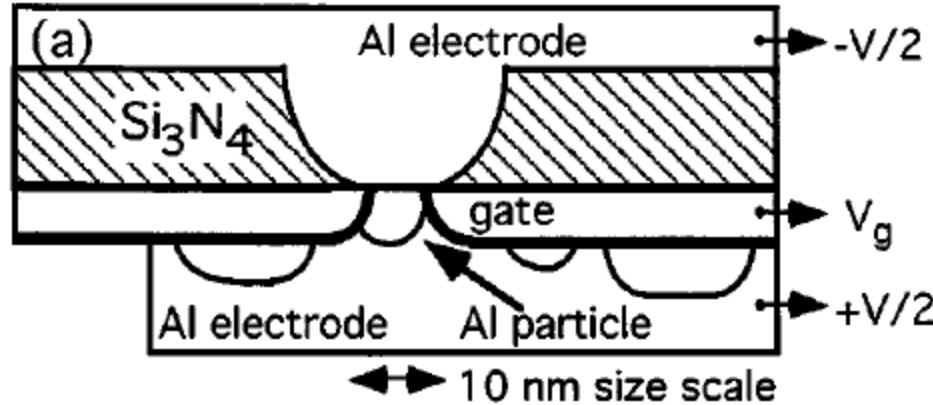
+Experimental control



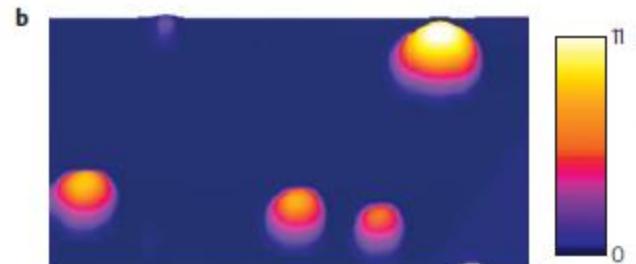
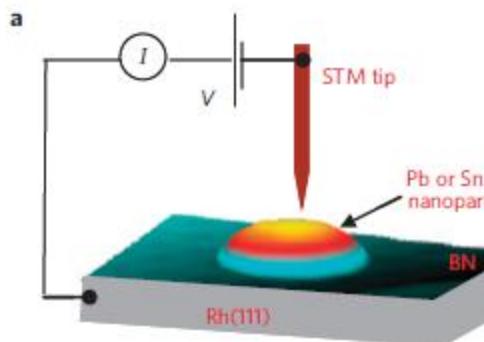
1966



+Predictive power



1995



Now

Next

Theory

Heterostructures
Collections of grains

Topology

Non-equilibrium

Experiments

Control on high T_c
heterostructures

Control on grains
arrangements

Substantial enhancement of T_c

THANKS!